

Electrons, phonons and photons

momentum:

$$p = \frac{h}{\lambda}$$

or

$$p = \hbar k$$

photons: λ is very large (400 - 700 nm) $\Rightarrow$ momentum is negligible
 phonons: λ is very small (few angstroms) $\Rightarrow$ large momentum

wavevector

$$k = \frac{2\pi}{\lambda}$$

photons: $k \approx 0$
 phonons: k is large

Energy

photons

$E \sim 1.7 - 3.1$ eV in the visible range $\Rightarrow$ large energy

phonons

$E \sim 63$ meV in silicon $\Rightarrow$ small energy

Generation and Recombination

Definitions

$G_{0,i}$
equilib. each process

G_0
total

G_i : generation rate [cm^{-3}/s]

R_i : recombination rate [cm^{-3}/s]

G, R : total

optical:

generation: generation of carriers (not light)

$$G_{0,\text{rad}} = g_{\text{rad}}(T) \quad [\text{cm}^{-3}/\text{s}] \quad \text{only depends on } T$$

recombination:

$$R_{0,\text{rad}} = r_{\text{rad}}(T) \cdot n_0 \cdot p_0$$

constant depending on
material and temperature
(how easy it is to
emit a photon)

very small in Si

$$r_{\text{rad}} = 2 \times 10^{-15} \text{ cm}^3/\text{s}$$

in GaAs at 300K:

$$r_{\text{rad}} = 7 \times 10^{-10} \text{ cm}^3/\text{s}$$

In TE:

$$G_{0,\text{rad}} = R_{0,\text{rad}} \Rightarrow g_{\text{rad}}(T) = r_{\text{rad}}(T) \cdot n_0 \cdot p_0 = r_{\text{rad}}(T) \cdot n_i^2$$

For example:

Consider n-type Si with $N_D = 10^{17} \text{ cm}^{-3}$ at R.T.

What are the generation and recombination rates due to band-to-band optical processes?

$$\text{In TE: } G_{0,\text{rad}} = R_{0,\text{rad}}$$

$$\text{for } N_D = 10^{17} \text{ cm}^{-3} \Rightarrow n_0 = 10^{17} \text{ cm}^{-3} \quad \text{thus } p_0 \approx 10^3 \text{ cm}^{-3}$$

$$\left. \begin{aligned} n_0 \cdot p_0 &= n_i^2 \\ n_i &\approx 10^{10} \text{ cm}^{-3} \end{aligned} \right\}$$

$$\therefore G_{0,\text{rad}} = R_{0,\text{rad}} = r_{\text{rad}} \cdot n_0 \cdot p_0 =$$

$$= 2 \times 10^{-15} \frac{\text{cm}^3}{\text{s}} \cdot 10^{17} \text{ cm}^{-3} \cdot 10^3 \text{ cm}^{-3} =$$

$$= 2 \times 10^5 \text{ cm}^{-3}/\text{s}$$

Out of equilibrium situation:

Optical generation rate: still the same

$$G_{\text{rad}} = G_{\text{rad}} = r_{\text{rad}} \cdot n_0 \cdot p_0$$

Optical recombination rate:

$$R_{\text{rad}} = r_{\text{rad}} \cdot n \cdot p$$

not in equilibrium

Net recombination rate:

$$U_{\text{rad}} = R_{\text{rad}} - G_{\text{rad}} = r_{\text{rad}} (np - \underbrace{n_0 p_0}_{n_i^2})$$

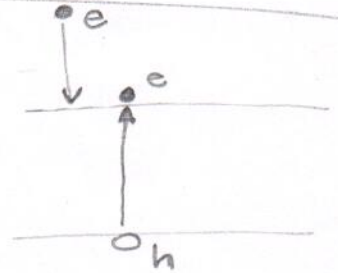
$$\left\{ \begin{array}{l} \text{if: } np > n_i^2 \Rightarrow \text{recombination} \\ np < n_i^2 \Rightarrow \text{generation} \end{array} \right.$$

Auger recombination:

Generation: $G_{0, \text{eeh}} = g_{\text{eeh}}(T) \cdot n_0$

equil. electron-electron-hole

generation:



proportional to
number of hot electrons,
which is related to number of electrons

recombination: $R_{0, \text{eeh}} = r_{\text{eeh}} \cdot n_0^2 \cdot p_0$

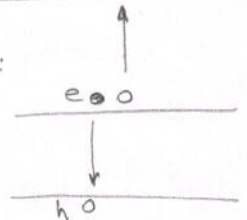
constant depending
on temperature
and
material

in TE: $g_{\text{eeh}} = r_{\text{eeh}} \cdot n_0 p_0$

For Si at R.T.:

$$r_{\text{eeh}} = 1.8 \times 10^{-31} \text{ cm}^6/\text{s}$$

recombination:



For example: n-type Si, $N_D = 10^{17} \text{ cm}^{-3}$ at room temperature

r_{eeh} for Si is $1.8 \times 10^{-31} \text{ cm}^6/\text{s}$

$$G_{0, \text{eeh}} = R_{0, \text{eeh}} = r_{\text{eeh}} \cdot n_0^2 p_0 = 1.8 \times 10^{-31} \frac{\text{cm}^6}{\text{s}} \cdot 10^{17} \text{ cm}^{-3} \cdot 10^3 \text{ cm}^{-3} = 1.8 \times 10^6 \text{ cm}^{-3}/\text{s}$$

out of equilibrium:

$$\left. \begin{aligned} G_{ech} &= g_{ech} \cdot n \\ R_{ech} &= r_{ech} \cdot n^2 \cdot p \end{aligned} \right\} U_{ech} = R_{ech} - G_{ech} = r_{ech} \cdot n (np - n_i^2)$$

Trap - assisted:

* Derivation is uploaded on moodle.

ⓐ exercise to demonstrate this

for $E_t = E_i$ with N_t

Probability of occupation: $f(E_t) = f(E_i) = \frac{1}{1 + \exp\left(\frac{E_i - E_F}{kT}\right)} = \frac{n_i}{n_i + p_0}$

occupied traps: $N_t \times f(E_i) = n_{to}$

unoccupied traps: $N_t - (N_t f(E_i)) = (N_t - n_{to})$

Rates:

capture of electrons: $r_{o,ec} = c_e \cdot n_0 \cdot (N_t - n_{to})$

electron
capture
coefficient

electrons

available traps

electron emission:

$$r_{o,ee} = e_e \cdot n_{to}$$

electron
emission
coefficient

occupied traps

* conduction band is largely empty of electrons

hole capture:

$$r_{o,hc} = c_h \cdot p_0 \cdot n_{to}$$

hole
capture

holes

occupied traps

hole emission:

$$r_{o,he} = e_h \cdot (N_t - n_{to})$$

* valence band has enough electrons.

for n-type semiconductor:

$$n_i \gg p_0$$

$$\Rightarrow r_{0,ec} = r_{0,ee} \approx \frac{n_i}{\tau_{eo}}$$

$$r_{0,hc} = r_{0,he} \approx \frac{p_0}{\tau_{ho}}$$

$$\text{if } \tau_{eo} \text{ not too different from } \tau_{ho} \Rightarrow r_{0,ee} \gg r_{0,hc}$$

• traps communicate more with CB than with VB

Exercise: n-type SC with $N_D = 10^{17} \text{ cm}^{-3}$ at room temperature

• There is a dominant trap at E_i with $N_t = 10^{16} \text{ cm}^{-3}$

• For this trap: $c_e = c_h = 10^{-11} \text{ cm}^3/\text{s}$

② Rate at which trap-assisted generation and recombination takes place?

③ Occupation probability of trap?

$$\tau_{eo} = \frac{1}{N_t c_e} = \frac{1}{10^{16} \times 10^{-11}} = 10^{-5} \text{ s}$$

$$\tau_{ho} = \frac{1}{N_t c_h} = \frac{1}{10^{16} \times 10^{-11}} = 10^{-5} \text{ s}$$

• rate of exchange of electrons between traps and CV:

$$r_{0,ce} = r_{0,ec} \approx \frac{n_i}{\tau_{eo}} = \frac{10^{10} \text{ cm}^{-3}}{10^{-5} \text{ s}} = 10^{15} \text{ cm}^{-3}/\text{s}$$

• for holes

$$r_{0,ch} = r_{0,eh} = \frac{p_0}{\tau_{ho}} = \frac{10^3}{10^{-5}} = 10^8 \text{ cm}^{-3}/\text{s}$$

Even though the trap is at E_i , $c_e = c_h$, trap "communicates" 7-orders of magnitude more with CV. $\Rightarrow$ because it is a n-type material.

$$b) f(E_c) = f(E_i) = \frac{n_i}{n_i + p_0} = \frac{10^{10}}{10^{10} + 10^3} \approx 1$$

E_F is well above E_i .

Very useful exercise:

calculate n_i for Si at room temperature:

$$n_i = \sqrt{N_c \cdot N_v} \cdot \exp\left(\frac{-E_g}{2kT}\right)$$

$$= \sqrt{2.86 \times 10^{19} \text{ cm}^{-3} \times 3.16 \times 10^{19} \text{ cm}^{-3}} \cdot \exp\left(-\frac{1.124 \text{ eV}}{2.259 \times 10^{-3} \text{ eV}}\right) = \underline{1.12 \times 10^{10} \text{ cm}^{-3}}$$

Useful to keep this value in mind

$$\boxed{n_i = 1.12 \times 10^{10} \text{ cm}^{-3}} \approx 10^{10} \text{ cm}^{-3} \text{ for quick calculations}$$

Proof 2 :

$$e_e = c_e \cdot n_o \cdot \frac{N_t - n_{to}}{n_{to}}$$

but we know :

$$n_i^2 = n_o \cdot p_o$$

$$n_{to} = N_t \cdot \frac{n_i}{n_i + p_o}$$

So :

$$e_e = c_e \cdot \frac{n_i^2}{p_o} \cdot \frac{\cancel{N_t} \left(1 - \frac{n_i}{n_i + p_o} \right)}{\cancel{N_t} \cdot \left(\frac{n_i}{n_i + p_o} \right)} =$$

$$= c_e \cdot \frac{n_i^2}{p_o} \cdot \frac{\frac{p_o}{n_i + p_o}}{\frac{n_i}{n_i + p_o}} = c_e \cdot n_{i/}$$

Recombination and Generation:

$$T.E \Rightarrow P.n = n_i^2$$

out of equilibrium $\Rightarrow$

$$P.n > n_i^2 \Rightarrow \text{recombination}$$
$$P.n < n_i^2 \Rightarrow \text{generation}$$

Recombination rate: $R = r_{\text{rad}} \cdot P.n$

$$r_{\text{rad}} \approx 10^{-10} \text{ cm}^3/\text{s} \quad \text{direct band gap}$$
$$\approx 10^{-15} \text{ cm}^3/\text{s} \quad \text{indirect band gap (unlikely)}$$

Net transition rate:

$$U = R - G = r_{\text{rad}} (P.n - n_i^2)$$

Under low injection: ($\Delta p' \ll N_D$)

n-type semiconductor:

$$n \approx N_D \Rightarrow U = r_{\text{rad}} \left(\underbrace{n \cdot p}_{n_i^2} + n \cdot p' - n_i^2 \right)$$

$$p = p_0 + \Delta p'$$

$$U \approx r_{\text{rad}} \cdot N_D \cdot p' \equiv \frac{p'}{\tau_p} \quad \begin{array}{l} \text{extra} \\ \text{\# carriers} \\ \text{per} \\ \text{time.} \end{array}$$

Definition of lifetime:

Average time for a minority carrier to recombine:

$$\left\{ \begin{array}{l} \tau_p = \frac{1}{r_{\text{rad}} \cdot N_D} : \text{n-type} \\ \tau_n = \frac{1}{r_{\text{rad}} \cdot N_A} : \text{p-type} \end{array} \right.$$

Exercise:

Consider n-type Si, under low level approximation, and $N_D = 10^{17} \text{ cm}^{-3}$

Assume a trap at E_t with $N_t = 10^{16} \text{ cm}^{-3}$ and $C_e = C_h = 10^{-11} \text{ cm}^3/\text{s}$

$$\Rightarrow N_0 \sim N_D = 10^{17} \text{ cm}^{-3}$$

We know that:

$$\tau_{\text{rad}} = \frac{1}{\Gamma_{\text{rad}} \cdot n_0} = \frac{1}{2 \times 10^{-15} \text{ cm}^3/\text{s} \times 10^{17} \text{ cm}^{-3}} = 5 \text{ ns}$$

$$\tau_{\text{Auger}} = \frac{1}{K_{\text{ech}} \cdot n_0^2} = \frac{1}{1.8 \times 10^{-31} \text{ cm}^6/\text{s} \times (10^{17} \text{ cm}^{-3})^2} = 0.56 \text{ ns}$$

$$\tau_{\text{tr}} = \tau_{\text{ho}} = \frac{1}{10^{16} \text{ cm}^{-3} \times 10^{-11} \text{ cm}^3/\text{s}} = 10 \mu\text{s}$$

thus

$$\tau = \frac{1}{\frac{1}{\tau_{\text{rad}}} + \frac{1}{\tau_{\text{Auger}}} + \frac{1}{\tau_{\text{tr}}}} =$$

$$= \frac{1}{\frac{1}{5 \times 10^{-3}} + \frac{1}{5.6 \times 10^{-4}} + \frac{1}{10^{-5}}} = 9.8 \mu\text{s}$$

Recombination lifetime

$\Rightarrow$ is dominated by

trap-assisted

recombinations.

$$\frac{dn}{dt} = \frac{dp}{dt} = G - R$$

if $G > R \Rightarrow \frac{dn}{dt} > 0, \frac{dp}{dt} > 0 \Rightarrow$ we generate carriers n and p

with external generation (for example shining light):

$$G = G_{\text{ext}} + G_{\text{int}}$$

then:

$$\frac{dn}{dt} = \frac{dp}{dt} = G_{\text{ext}} - \underbrace{G_{\text{int}} + R}_U = G_{\text{ext}} - U$$

Under low-level injection:

$$U_i \approx \frac{n'}{\tau_i}$$

$$\frac{dn}{dt} = \frac{dn'}{dt}$$

Since $n = n_0 + n'$ and assuming that n_0 is time independent.

$$\boxed{\frac{dn'}{dt} = G_{\text{ext}} - \frac{n'}{\tau}}$$

Case 1: static case where $\frac{dn}{dt} = 0$ (nothing changes as a function of time)

$$\Rightarrow n' = G_{\text{ext}} \cdot \tau$$

excess carriers is the generation rate times the carrier lifetime.

Case 2: if $G_{\text{ext}} = 0$

the homogeneous solution is $n' \propto e^{-t/\tau}$: excess carriers decay with a time τ

1) Turn-on transient :

• Lights are turned ON at $t=0$:

$$\begin{cases} t=0, n'=0 \\ t=\infty, n'=g_e \tau \end{cases}$$

• τ is a characteristic time of the problem

$$\Rightarrow n'(t) = g_e \tau (1 - e^{-t/\tau})$$

• $t \gg \tau$: rate of generation of carriers decrease towards 0

1. Initially : n' increases since $G > R$

2. As n' builds up, R increases but G does not change.

3. n' is constant when $G = R \Rightarrow$ steady-state is reached!

2) turn-off transient

• Lights are ON for $t < 0$

• $t > 0$, $G_{ext} = 0$

boundary conditions:

$$\begin{cases} t=0, n'=g_e \tau \\ t=\infty, n'=0 \text{ (lights are OFF)} \end{cases}$$

$$\Rightarrow n'(t) = g_e \tau e^{-t/\tau}$$

1. After equilibrium is disrupted, it takes $t > \tau$ to re-establish it.

2. After equilibrium is disrupted : R not balanced by $(G_{ext} + G_{int})$

$\Rightarrow R > G$, thus n' is reducing and so R

until $R = G$